

## Voltage tunable nonreciprocity of microwaves transmission with varactor-loaded planar meta-structures

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**Abstract** – We suggest new planar meta-structures “ferrite plate – varactor-loaded resonant element” that provide voltage tunable nonreciprocity of microwaves transmission in contrast to traditional magnetic control. The effect is observed by tuning varactor capacitance with changing reverse-bias voltages under the ferromagnetic resonance excitation near the varactor-loaded element resonance frequency. Wide tuning range has been obtained with nonreciprocity 19 dB between 5.4 and 5.9 GHz. Presented meta-structures are useful for development of quick-tunable nonreciprocal systems.

### I. INTRODUCTION

Last time of special interest is the study of voltage tunable varactor-loaded metamaterials. Many publications along with [1] are devoted to voltage tuning through the voltage controlled varactor capacitance. Besides, unique effects are observed, for example, nonlinear magnetoelectric coupling [2] and nonlinear optical activity [3]; there are no natural materials with such properties. In this presentation for the first time we demonstrate voltage tunable nonreciprocal resonance response at microwaves in contrast to traditional control methods by magnetic field. Previously we investigated nonreciprocal microwave resonance response and non-additive properties with complex planar meta-structures containing ferrite plate and grating of resonant elements forming surface plasmon-polaritons. In papers [4] and [5] it was shown that such planar meta-structures exhibit giant nonreciprocity in the transmission at the ferromagnetic resonance (FMR) frequencies at certain values of the static magnetic field  $H$  under conditions of a mutual influence between the FMR and grating resonance while nonreciprocal resonance effects are absent with free ferrite. The nonreciprocal FMR is due to the interaction between precessing spins in ferrite and a local magnetic field  $h$  of the surface wave with elliptic (circular) polarization formed by a grating (chains or single resonant element). The nonreciprocity  $\delta$  is the difference between transmission coefficients  $T$  corresponding to the opposite directions of magnetization  $H+$  and  $H-$  (opposite senses of spin precession) or is the difference between transmission coefficients  $T$  for propagating modes in the opposite directions.

Here we demonstrate in a rectangular waveguide magnetic and electric control of transmission nonreciprocity by both field  $H$  and small direct-current voltages  $V_{DC}$  with “ferrite plate - varactor loaded split-ring” meta-structure. Wide tuning range has been obtained also with varactor-loaded split-butterfly dipole.

### II. INVESTIGATED META-STRUCTURES.

#### EVOLUTION OF RESONANCE RESPONSE WITH STATIC MAGNETIC AND ELECTRIC FIELDS

Investigated meta-structure contains a plate of isotropic yttrium-iron garnet  $3Y_2O_3 \cdot 5Fe_2O_3$  (30 x 20 x 1.4 mm) and copper twice split ring loaded with voltage dependent nonlinear capacitance as varactor BB857 (Fig. 1). The varactor capacitance can be tuned between 6.5pF and 0.65 pF by changing reverse-bias voltages  $V_{DC}$  from 1 V to 30 V. Meta-structure is arranged along the rectangular waveguide axis in a transverse external static magnetic field  $H$  and oriented parallel to the narrow side wall. Evolution of amplitude-frequency resonance response (transmission coefficient  $T$ ) is investigated in dependence on the magnitude of applied field  $H$  under the magnetization reversal and in dependence on reverse-bias voltages  $V_{DC}$ .

Applying  $H$  field the **FMR** is excited and shifts to higher frequencies with  $H$  field increase.

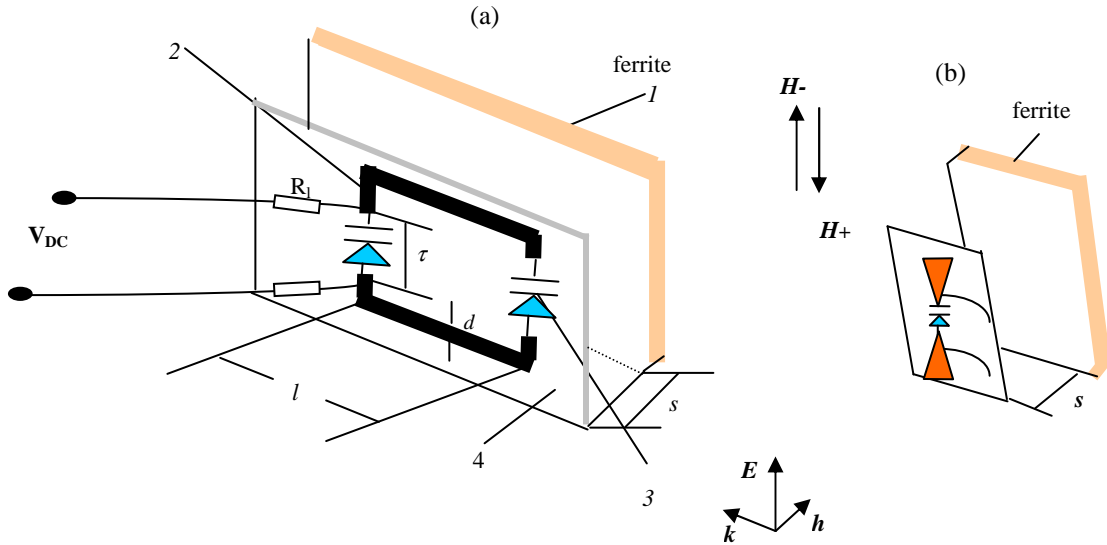


Fig. 1. Meta-structures: (a) “ferrite plate 1– varactors 3-loaded twice split-ring 2”, substrate 4,  $R_L = 100 \text{ k}\Omega$ ,  $l=18 \text{ mm}$ ,  $\tau=1.5 \text{ mm}$ ,  $d = 3 \text{ mm}$ ; (b) “ferrite plate – varactor-loaded split-dipole butterfly”.

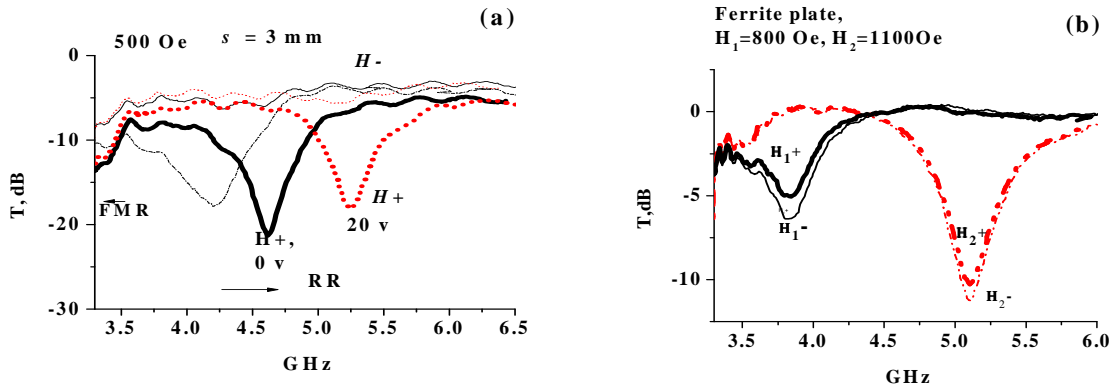


Fig.2. Frequency dependences of the transmission coefficient  $T$  in a rectangular waveguide, curves for  $H+$  (bold) and  $H-$  (thin) correspond to different directions of magnetization:  
(a) metastructure ( $s = 3 \text{ mm}$ ),  $H=0$ ,  $V_{DC} = 0$  (dash-dot);  $H = 500 \text{ Oe}$ ,  $V_{DC} = 0$  (solid),  $H = 500 \text{ Oe}$ ,  $V_{DC} = 20 \text{ V}$  (dot);  
(b) ferrite without ring, nonreciprocity is absent.

In Figure 2 we see evolution of resonance response with meta-structure (distance  $s = 3 \text{ mm}$ ) applying a static magnetic field  $H = 500 \text{ Oe}$  and voltage  $V_{DC} = 20 \text{ V}$ . When  $H = 0$ ,  $V_{DC} = 0$  the **FMR** is not excited but we see resonance response at  $4.2 \text{ GHz}$  (dash-dot curve) due to varactor-loaded split-ring resonance (**RR**). Under  $H = 500 \text{ Oe}$ , when the **FMR** is excited at about  $3 \text{ GHz}$ , the **RR** is observed at  $4.6 \text{ GHz}$ , the **RR** frequency shift is  $\Delta f_H = 0.4 \text{ GHz}$  with changing value of the  $H$  field from  $0$  to  $500 \text{ Oe}$ . Besides, nonreciprocity appears at the **RR** frequencies, level of resonant curve under the  $H-$  (thin curve) and under  $H+$  field (bold curve) is not the same, nonreciprocity  $\delta = T(H+) - T(H-)$  is  $15 \text{ dB}$ . Nonreciprocal **RR** can be tuned by changing the biasing voltage  $V_{DC}$ . Changing the  $V_{DC}$  from  $0$  to  $20 \text{ V}$  the **RR** frequency shifts to  $5.25 \text{ GHz}$  and the **RR** frequency

shift  $\Delta f_{DC} = 0.65$  GHz. The field  $H$  is much less than the field (1150 Oe) necessary to the FMR excitation at RR frequency.

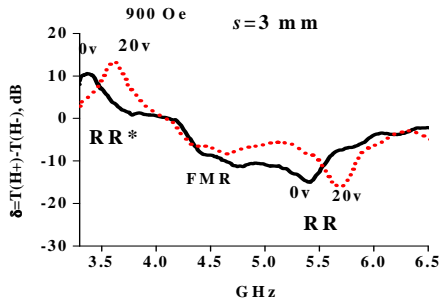


Fig. 3. Bifurcation. Evolution of frequency dependence of the nonreciprocity  $\delta$  with changing  $V_{DC}$  under  $s = 3$  mm:  $V_{DC}=0$  V (bold);  $V_{DC}=20$  V (dot)

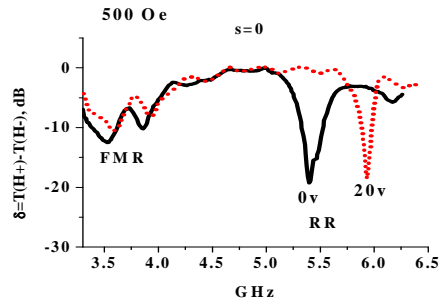


Fig. 4. Evolution of frequency dependence of the nonreciprocity  $\delta$  with changing  $V_{DC}$  under  $s=0$ :  $V_{DC}=0$  V (bold);  $V_{DC}=20$  V (dot)

With further increase of the field  $H$  when moving FMR approaches the RR first position one can observe bifurcation effect. In this case the second ring resonance ( $RR^*$ ) is excited lower than the first position and we observe two separate nonreciprocal resonances ( $RR_B$  and  $RR^*_B$ ) with opposite signs of the nonreciprocity  $\delta$  at different frequencies. In Fig. 3 we see nonreciprocity  $\delta$  corresponding to bifurcation situation under the field  $H = 900$  Oe. In this case the  $RR_B$  is excited at 5.4 GHz and  $RR^*_B$  is excited at 3.4 GHz. Two resonances can be tuned simultaneously by changing varactor biasing voltage. By changing the  $V_{DC}$  from 0 to 20 V the shift  $\Delta f_{DC}(RR_B) = 0.29$  GHz and the shift  $\Delta f^*_{DC}(RR^*_B) = 0.27$  GHz. In this time transmission at the FMR domain (around 4.5 GHz) is nonreciprocal but the FMR frequency is not practically tuned by changing voltages  $V_{DC}$ .

With decreasing the distance  $s$  coupling between ferrite and ring increases and the shift  $\Delta f_H$  increases. In this case the nonreciprocity  $\delta$  also increases. Fig.4 corresponds to  $s = 0$ . The shift  $\Delta f_H = 1.1$  GHz with changing  $H$  from 0 to 500 Oe; the nonreciprocity  $\delta = 19$  dB and the shift  $\Delta f_{DC} = 0.53$  GHz by changing  $V_{DC}$  from 0 to 20 V.

### III. CONCLUSION

So, the electrical tunability of nonreciprocity of microwaves transmission by tuning varactor capacitance has been demonstrated in waveguide with meta-structure containing a transversely magnetized ferrite and a varactor-loaded split-ring under interaction between precessing ferrite's spins and a magnetic field of the wave formed by the ring. Wide tuning range has been obtained also with varactor-loaded split-butterfly dipole.

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